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Generalization of Titchmarsh's Theorem for the Dunkl Transform in the Space $L^p_{\alpha}(\mathbb{R})$

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Abstract. In this paper, using a generalized Dunkl translation operator, we obtain a generalization of Titchmarsh's Theorem for the Dunkl transform for functions satisfying the (ψ, p) -Lipschitz Dunkl condition in the space $L_{p,\alpha} = L^p(\mathbb{R}, |x|^{2\alpha+1}dx)$, where $\alpha > -\frac{1}{2}$.

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1. Introduction and preliminaries

Dunkl operators are differential-difference operators introduced in 1989, by Dunkl [2]. On the real line, these operators, which are denoted by D_{α} , depend on a real parameter $\alpha > -\frac{1}{2}$.

In [1], we proved an analog of Titchmarsh's theorem for the Dunkl transform in the space $L_{2,\alpha}$. In this paper we prove a generalization of this theorem in the space $L_{p,\alpha}$, where 1 . For this purpose, we use a generalized Dunkl translation operator.

 $L_{p,\alpha} = L^p(\mathbb{R}, |x|^{2\alpha+1}dx); 1 , is the Banach space of measurable functions <math>f(t)$ on \mathbb{R} with finite norm

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$$||f||_{p,\alpha} = \left(\int_{-\infty}^{\infty} |f(x)|^p |x|^{2\alpha+1} dx\right)^{1/p}.$$

The Dunkl operator is a differential-difference operator D_{α}

$$D_{\alpha} = \frac{df(x)}{dx} + (\alpha + \frac{1}{2})\frac{f(x) - f(-x)}{x}, \ \alpha > -\frac{1}{2},$$

where $f \in L_{p,\alpha}$.

Let $j_{\alpha}(t)$ is a normalized Bessel function of the first kind

$$j_{\alpha}(t) = \frac{2^{\alpha} \Gamma(\alpha + 1) J_{\alpha}(t)}{t^{\alpha}},$$

where $J_{\alpha}(t)$ is a Bessel function of the first kind. The function $j_{\alpha}(t)$ is infinitely differentiable and even.

The Dunkl kernel defined by

$$e_{\alpha}(x) = j_{\alpha}(x) + ic_{\alpha}j_{\alpha+1}(x)$$

where $c_{\alpha} = (2\alpha + 2)^{-1}$.

Using the correlation

$$j'_{\alpha}(x) = -\frac{xj_{\alpha+1}(x)}{2(\alpha+1)}.$$

We have

$$e_{\alpha}(x) = j_{\alpha}(x) - ij'_{\alpha}(x). \tag{1}$$

The Dunkl transform is defined by

$$\widehat{f}(\lambda) = \int_{-\infty}^{\infty} f(x)e_{\alpha}(\lambda x)|x|^{2\alpha+1}dx, \ \lambda \in \mathbb{R}.$$

The inverse Dunkl transform is defined by the formula

$$f(x) = (2^{\alpha+1}\Gamma(\alpha+1))^{-2} \int_{-\infty}^{\infty} \widehat{f}(\lambda)e_{\alpha}(-\lambda x)|\lambda|^{2\alpha+1}d\lambda.$$

Plancherel's theorem and the Marcinkiewics interpolation theorem (see [3]) we get for $f \in L_{p,\alpha}$ with 1 and <math>q such that $\frac{1}{p} + \frac{1}{q} = 1$,

$$\|\widehat{f}\|_{q,\alpha} \leqslant C\|f\|_{p,\alpha},\tag{2}$$

where C is a positive constant.

K. Trimèche has introduced in [4] the generalized Dunkl translation operator τ_h , $h \in \mathbb{R}$, we have

$$\widehat{(\tau_h f)}(x) = e_{\alpha}(xh)\widehat{f}(x). \tag{3}$$

The function $j_{\alpha}(x)$ is defined also by

$$j_{\alpha}(z) = \Gamma(\alpha + 1) \sum_{n=0}^{\infty} \frac{(-1)^n (z/2)^{2n}}{n! \Gamma(n+\alpha+1)}, \ z \in \mathbb{C}.$$

$$\tag{4}$$

Moreover, from (4) we see that

$$\lim_{z \to 0} \frac{j_{\alpha}(z) - 1}{z^2} \neq 0$$

by consequence, there exist c > 0 and $\eta > 0$ satisfying

$$|z| \leqslant \eta \Longrightarrow |j_{\alpha}(z) - 1| \geqslant c|z|^2.$$
 (5)

2. Main Result

In this section we give the main result of this paper. We need first to define (ψ, p) -Lipschitz Dunkl class.

Definition 2.1 A function $f \in L_{p,\alpha}$ is said to be in the (ψ, p) -Lipschitz Dunkl class, denoted by $Lip(\psi, p)$, if

$$\|\tau_h f(x) + \tau_{-h} f(x) - 2f(x)\|_{p,\alpha} = O(\psi(h)) \text{ as } h \longrightarrow 0,$$

where $\psi(t)$ is a continuous increasing function on $[0,\infty)$, $\psi(0)=0$ and $\psi(ts)=\psi(t)\psi(s)$ for all $t,s\in[0,\infty)$.

Theorem 2.2 Let f(x) belong to $Lip(\psi, p)$. Then

$$\int_{|\lambda| \geqslant r} |\widehat{f}(\lambda)|^q |\lambda|^{2\alpha+1} d\lambda = O(\psi(r^{-q})) \text{ as } r \longrightarrow +\infty.$$

Proof Let $f \in Lip(\psi, p)$. Then we have

$$\|\tau_h f(x) + \tau_{-h} f(x) - 2f(x)\|_{p,\alpha} = O(\psi(h)) \text{ as } h \longrightarrow 0$$

From formulas (1) and (2), we have the Dunkl transform of $\tau_h f(x) + \tau_{-h} f(x) - 2f(x)$ is $2(j_{\alpha}(\lambda h) - 1)\widehat{f}(\lambda)$.

By (2), we obtain

$$\left(\int_{-\infty}^{\infty} 2^q |j_{\alpha}(\lambda h) - 1|^q |\widehat{f}(\lambda)|^q |\lambda|^{2\alpha + 1} d\lambda\right)^{1/q} \leqslant C \|\tau_h f(x) + \tau_{-h} f(x) - 2f(x)\|_{p,\alpha}$$

From (5), we have

$$\int_{\frac{\eta}{2h} \leqslant |\lambda| \leqslant \frac{\eta}{h}} |1 - j_{\alpha}(\lambda h)|^{q} |\widehat{f}(\lambda)|^{q} |\lambda|^{2\alpha + 1} d\lambda \geqslant \frac{c^{q} \eta^{2q}}{2^{2q}} \int_{\frac{\eta}{2h} \leqslant |\lambda| \leqslant \frac{\eta}{h}} |\widehat{f}(\lambda)|^{q} |\lambda|^{2\alpha + 1} d\lambda$$

There exists then a positives constants C_1 and K_1 such that

$$\int_{\frac{\eta}{2h} \leqslant |\lambda| \leqslant \frac{\eta}{h}} |\widehat{f}(\lambda)|^q |\lambda|^{2\alpha+1} d\lambda \leqslant C_1 \int_{-\infty}^{\infty} |1 - j_{\alpha}(\lambda h)|^q |\widehat{f}(\lambda)|^q |\lambda|^{2\alpha+1} d\lambda
\leqslant K_1 \psi^q(h) = K_1 \psi(h^q).$$

Then

$$\int_{r \leqslant |\lambda| \leqslant 2r} |\widehat{f}(\lambda)|^q |\lambda|^{2\alpha + 1} d\lambda \leqslant K\psi(r^{-q}).$$

where $K = K_1 \psi(\eta^q 2^{-q})$.

Of course

$$\int_{|\lambda| \geqslant r} |\widehat{f}(\lambda)|^{q} |\lambda|^{2\alpha+1} d\lambda = \left(\int_{r \leqslant |\lambda| \leqslant 2r} + \int_{2r \leqslant |\lambda| \leqslant 4r} + \int_{4r \leqslant |\lambda| \leqslant 8r} + \dots \right) |\widehat{f}(\lambda)|^{q} |\lambda|^{2\alpha+1} d\lambda
\leqslant K \psi(r^{-q}) + K \psi((2r)^{-q}) + K \psi((4r)^{-q}) + \dots
\leqslant K \psi(r^{-q}) + K \psi(2^{-q}) \psi(r^{-q}) + K \psi((2^{-q})^{2}) \psi(r^{-q}) + \dots
\leqslant K \psi(r^{-q}) (1 + \psi(2^{-q}) + \psi((2^{-q})^{2}) + \dots).$$

We have $\psi(2^{-q}) < 1$, then

$$\int_{|\lambda| \geqslant r} |\widehat{f}(\lambda)|^q |\lambda|^{2\alpha + 1} d\lambda \leqslant C_2 \psi(r^{-q}),$$

where $C_2 = K(1 - \psi(2^{-q}))^{-1}$.

Finally, we get

$$\int_{|\lambda| \geqslant r} |\widehat{f}(\lambda)|^q |\lambda|^{2\alpha+1} d\lambda = O(\psi(r^{-q})) \text{ as } r \longrightarrow \infty.$$

Thus, the proof is finished.

Conclusion 3.

In this work we have succeeded to generalise the theorem in [1] for the Dunkl transform in the space $L_{p,\alpha}$. We proved that f(x) belong to $Lip(\psi,p)$. Then

$$\int_{|\lambda|\geqslant r} |\widehat{f}(\lambda)|^q |\lambda|^{2\alpha+1} d\lambda = O(\psi(r^{-q})) \text{ as } r \longrightarrow +\infty.$$

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