

Higher rank Einstein solvmanifolds

M. Zarghani

Department of Mathematics, Tarbiat Modares University, Tehran, Iran

E-mail: zarghanim@yahoo.com

ABSTRACT. In this paper we study the structure of standard Einstein solvmanifolds of arbitrary rank. Also the validity of a variational method for finding standard Einstein solvmanifolds is proved.

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1. INTRODUCTION

General form of standard Einstein solvmanifolds were determined by Jense Heber (see [2]). Later, Gorge Lauret deeply studied this kind of manifolds. Solvable Lie group endowed with the left invariant Riemannian metric is called solvmanifold. Let S be a simply connected Lie group with the corresponding Lie algebra s endowed with the inner product determined by $\langle \cdot, \cdot \rangle$ and solvable Lie bracket $[\cdot, \cdot]$. We call S a higher rank solvmanifold if

$$s = n \oplus a; \quad n = [s, s], \quad a = n^\perp,$$

where n is a metric nilpotent Lie algebra of dimension k . The codimension n is called the rank of S . The solvable Lie group $(S, [\cdot, \cdot], \langle \cdot, \cdot \rangle)$ is called standard if a is abelian and it is said to be Einstein if its Ricci tensor $ric_{[\cdot, \cdot]}$ satisfies $ric \langle \cdot, \cdot \rangle = c \langle \cdot, \cdot \rangle$, for some $c \in \mathbb{R}$. s is called a metric solvable extension of n if the restriction of the Lie bracket and inner product of s to n coincide respectively with the Lie bracket and inner product of n . Let $\Lambda^2 n^* \otimes n$

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be the vector space of all bilinear skew-symmetric maps from $n \times n$ to n . There is a natural action of $GL(k)$ on $\Lambda^2 n^* \otimes n$ which is given by

$$\phi.\mu(X, Y) = \phi\mu(\phi^{-1}X, \phi^{-1}Y); \quad X, Y \in n, \quad \phi \in GL(k), \quad \mu \in \Lambda^2 n^* \otimes n.$$

Let N denote a simply connected nilpotent Lie group with Lie algebra (n, μ) endowed with the left invariant Riemannian metric $\langle \cdot, \cdot \rangle_n$, where μ is a nilpotent Lie algebra on n . The Ricci operator $R_\mu : n \rightarrow n$ of N is defined by

$$\langle R_\mu \cdot, \cdot \rangle_n = \text{ric} \langle \cdot, \cdot \rangle_n.$$

This operator is reduced to

$$(1.1) \quad \langle R_\mu X, Y \rangle = -\frac{1}{2} \sum_{i,j} \langle \mu(X, X_i), X_j \rangle \langle \mu(Y, X_i), X_j \rangle \\ + \frac{1}{4} \sum_{i,j} \langle \mu(X_i, X_j), X \rangle \langle \mu(X_i, X_j), Y \rangle,$$

for all $X, Y \in n$, where $\{X_1, X_2, \dots, X_k\}$ is any orthonormal basis of n . The inner product $\langle \cdot, \cdot \rangle_n$ determines an inner product on $\Lambda^2 n^* \otimes n$, denoted by $\langle \cdot, \cdot \rangle$ and given by

$$\langle \mu, \lambda \rangle = \sum_{i,j,k} \langle \mu(X_i, X_j), X_k \rangle \langle \lambda(X_i, X_j), X_k \rangle.$$

Also it naturally determines a norm on $\Lambda^2 n^* \otimes n$ defined by

$$\forall \lambda \in \Lambda^2 n^* \otimes n \quad \|\lambda\| = \sum_{ijk} \langle \lambda(X_i, X_j), X_v \rangle^2.$$

Consider the Riemannian function

$$F_k : \Lambda^2 n^* \otimes n \longrightarrow \mathbb{R}, \quad F_k(\mu) = \text{tr} R_\mu^2$$

and the sphere S_r given by

$$S_r = \{\mu \in \Lambda^2 n^* \otimes n; \|\mu\|^2 = 2r^2\},$$

for some $r \in \mathbb{R}$. Let \mathfrak{N}_k be the vector space of all nilpotent Lie brackets on n and $Der(\mu)$ be the Lie algebra of all derivations on n . Then $\mu \in \mathfrak{N}_k$ is called a Ricci soliton if $R_\mu = cI + D$, for some $D \in Der(\mu)$ and $c \in \mathbb{R}$.

In [6], Jorge Lauret has proved that the standard Einstein solvmanifolds are exactly the critical points of modified scalar curvature function $F_k|_{S_1}$.

Theorem 1.1. [6]. *For $\mu \in \mathfrak{N}_k \cap S_1$, the following statements are equivalent:*

- (i) μ is a critical point of $F_k|_{S_1}$.
- (ii) μ is a critical point of $F_k|_{GL(k).\mu \cap S_1}$.
- (iii) μ admits a rank-one extension which is Einstein.
- (iv) μ is a Ricci soliton.

2. EINSTEIN SOLVMANIFOLDS OF RANK ≥ 1

In this section, the structure of standard Einstein solvmanifolds is introduced. The next lemma provides some useful properties of solvable Lie bracket and inner product of a solvable Lie group. Then, we extend the Ricci soliton in [4] for any arbitrary rank. We call it multiple Ricci soliton.

Lemma 2.1. [2]. *Let $(s = n \oplus a, [\cdot, \cdot], \langle \cdot, \cdot \rangle)$ be a metric solvable extension of $(n, \mu, \langle \cdot, \cdot \rangle)$, where for every $0 \neq A \in a$, ad_A is nonzero and symmetric, then*

- (i) $\langle R_{[\cdot, \cdot]}A, B \rangle = -tr(ad_A ad_B)$, for all $A, B \in a$.
- (ii) $\langle R_{[\cdot, \cdot]}A, X \rangle = 0$, for all $A \in a, X \in n$.
- (iii) $R_{[\cdot, \cdot]}|_n = -ad_Z|_n + R_\mu$, where $\langle Z, X \rangle = tr(ad_X)$, for all $X \in n$.

Definition 2.2. $0 \neq \mu \in \mathfrak{N}_k$ is called a multiple Ricci soliton of degree r if

- (a) $R_\mu = c_\mu I + D_\mu$; $D_\mu \in Der(\mu)$, $c_\mu \in \mathbb{R}$.
- (b) There are nonzero symmetric derivations D_i , $1 \leq i \leq r$, such that

$$D_\mu = D_1 + D_2 + \dots + D_r, \quad tr D_i D_j = -\delta_{ij} c_\mu tr D_i.$$

Remark 2.3. If μ is a multiple Ricci soliton of degree r , then μ is a multiple Ricci soliton of degree less than r . Therefore, μ is the critical point of $F_k|_{S_r \cap GL(k) \cdot \mu}$.

Using Definition 2.1, we study the structure of standard Einstein solvmanifolds as follows.

Proposition 2.4. *For $0 \neq \mu \in \mathfrak{N}_k \cap S_r$, the following statements are equivalent:*

- (i) μ admits a metric extension which is Einstein.
- (ii) μ is a multiple Ricci soliton.

Proof. Let the Lie algebra (n, μ) admit an Einstein metric extension S with corresponding Lie algebra $(s = n \oplus a, [\cdot, \cdot], \langle \cdot, \cdot \rangle)$ such that $dim(a) = r$. Let $\{H_1, H_2, \dots, H_r\}$ be an orthonormal basis for a and Z be the mean curvature vector field for the simply connected Lie group N with Lie algebra n . A straightforward calculation shows that $D_\mu = D_1 + D_2 + \dots + D_r$, where $D_\mu := ad_Z|_n$ and $D_i = tr(ad_{H_i})ad_{H_i}|_n$. $[\cdot, \cdot]$ is the Lie bracket. Hence, $D_i \mu(\cdot, \cdot) = \mu(\cdot, D_i \cdot) + \mu(D_i \cdot, \cdot)$; that is to say, D_i 's are derivations on n . Suppose that D_i 's and D_μ are symmetric (see [2; 4.10]). Let $Z_i = tr(ad_{H_i})H_i$, then Lemma 2.1 implies that

$$tr D_i D_j = tr(ad_{Z_i} ad_{Z_j}) = -\langle R_{[\cdot, \cdot]}Z_i, Z_j \rangle = -c_\mu \langle Z_i, Z_j \rangle = -\delta_{ij} c_\mu tr D_i.$$

Also $R_\mu = c_\mu I + D_\mu$, for some $c_\mu \in \mathbb{R}$. Therefore, μ is a multiple Ricci soliton.

Conversely, let μ be a multiple Ricci soliton i.e.

- (a) $R_\mu = c_\mu I + D_\mu$; $D_\mu \in Der(n, \mu)$, $c_\mu \in \mathbb{R}$.

(b) There are symmetric derivations D_i , $1 \leq i \leq r$, such that

$$D_\mu = D_1 + D_2 + \dots + D_r, \quad \text{tr} D_i D_j = -\delta_{ij} c_\mu \text{tr} D_i.$$

Let $(n, \mu, \langle \cdot, \cdot \rangle_n)$ be a Lie algebra with orthonormal basis $\{X_1, X_2, \dots, X_k\}$. We define Lie algebra s with a simply connected Lie group S as follows

$$s = n \oplus \sum_i \mathbb{R} Z_i,$$

endowed with the inner product $\langle \cdot, \cdot \rangle$ defined by

$$\langle Z_i, Z_j \rangle = \delta_{ij} \text{tr} D_i, \quad \langle Z_i, n \rangle = 0, \quad \langle \cdot, \cdot \rangle|_{n \times n} = \langle \cdot, \cdot \rangle_n.$$

Also, Lie bracket $[\cdot, \cdot]$ is defined by

$$[Z_i, Z_j] = 0, \quad [Z_i, X_j] = -[X_j, Z_i] = D_i X_j, \quad [\cdot, \cdot]|_{n \times n} = \mu.$$

Clearly $[\cdot, \cdot]$ is a Lie bracket, since D_i 's are derivations. $\{D_\mu X_1, D_\mu X_2, \dots, D_\mu X_k\}$ is a linearly independent set which generates a subalgebra of $[s, s]$. Therefore, $n = [s, s]$. μ is nilpotent hence $[\cdot, \cdot]$ is a solvable Lie bracket. Finally using Lemma 2.1, we have

$$\begin{aligned} \langle R_{[\cdot, \cdot]} Z_i, Z_j \rangle &= -\text{tr}(D_i D_j) = \delta_{ij} c_\mu \text{tr} D_i = c_\mu \langle Z_i, Z_j \rangle, \quad \langle R_{[\cdot, \cdot]} Z_i, n \rangle = 0, \\ \langle R_{[\cdot, \cdot]} X_i, X_j \rangle &= \langle (-D_\mu + R_\mu) X_i, X_j \rangle = \langle c_\mu X_i, X_j \rangle = c_\mu \langle X_i, X_j \rangle, \end{aligned}$$

which implies that $\langle \cdot, \cdot \rangle_s$ is a Einstein metric. This completes the proof. \square

Using Proposition 2.1, we get a higher rank Einstein solvmanifold as the direct sum of the Lie algebras.

Proposition 2.5. *If nonzero nilpotent Lie brackets μ_1 and μ_2 are Ricci solitons, then $\mu = \mu_1 \oplus \mu_2$ is a multiple Ricci soliton of degree 2.*

Proof. μ_1 and μ_2 are Ricci solitons i.e.

$$(2.1) \quad R_{\mu_i} = c_{\mu_i} I + D_{\mu_i}; \quad D_{\mu_i} \in \text{Der}(\mu_i), \quad c_{\mu_i} \in \mathbb{R}, \quad i = 1, 2.$$

Up to isometry and scaling we can determine norms of μ_1 and μ_2 such that $c_{\mu_1} = c_{\mu_2}$. Set

$$D_\mu = \begin{bmatrix} D_{\mu_1} & 0 \\ 0 & D_{\mu_2} \end{bmatrix},$$

$$R_\mu = \begin{bmatrix} R_{\mu_1} & 0 \\ 0 & R_{\mu_2} \end{bmatrix},$$

Then $R_\mu = c_{\mu_1} I + D_\mu$, $R_\mu = R_{\mu_1} \oplus R_{\mu_2}$ and $D_\mu = D_{\mu_1} \oplus D_{\mu_2}$. Also, by Theorem 1.3, $\text{tr} D_{\mu_i} D_{\mu_j} = -\delta_{ij} c_{\mu_1} \text{tr} D_{\mu_i}$; $i = 1, 2$. Therefore μ is a multiple Ricci soliton which admits a 2-rank Einstein solvable extension. \square

Corollary 2.6. *If nonzero nilpotent Lie brackets μ_i 's, $1 \leq i \leq r$, are Ricci solitons, then $\mu = \mu_1 \oplus \mu_2 \oplus \dots \oplus \mu_r$ is a multiple Ricci soliton of degree r which admits an Einstein solvable extension of rank r .*

Remark 2.7. There exist 31 Ricci soliton nonzero Lie algebras of dimension 6 (see [7]), which by direct sum of them, we can obtain a lot of multiple Ricci soliton nilpotent Lie algebras.

3. STANDARD METHODS

The goal of this section is to present certain results from [4] and [6]. In view of [4], Jorge Lauret has used a variational method for finding standard Einstein solvmanifolds. We will demonstrate this method in Theorem 3.1. We first give some preliminaries.

Lemma 3.1. (*Lagrange multiplier theorem*) [1]. *Let P and M be smooth manifolds and $g : M \rightarrow P$ be a smooth submersion. Let $f : M \rightarrow \mathbb{R}$ be C^r , $m \in M$ and $p \in P$ such that $m \in g^{-1}(p)$, then the following statements are equivalent:*

- (i) m is a critical point of $f|_{g^{-1}(p)}$.
- (ii) There are $\lambda \in T_p^*M$ such that $T_m f = \lambda \circ T_m g$.

The vector space \mathfrak{N}_k is $GL(k)$ -invariant, so we can refine Theorem 1.1 more accurately as follows.

Lemma 3.2. *For $0 \neq \mu \in \mathfrak{N}_k$ and $\psi \in GL(k)$ the following statements are equivalent:*

- (i) $\psi.\mu$ is a Ricci soliton.
- (ii) $\psi.\mu$ is a critical point of $F_k|_{S_r \cap GL(k).\mu}$.
- (iii) ψ is a solution of the following system of equations:

$$\begin{cases} \|\phi.\mu\|^2 = 2r^2 \\ \frac{\partial F_k(\phi.\mu)}{\partial \phi_{ij}} = t \frac{\partial(\|\phi.\mu\|)}{\partial \phi_{ij}} \end{cases}$$

where $t \in \mathbb{R}$ and $\phi \in GL(k)$.

Proof. Let $g(\lambda) = \frac{1}{2}\|\lambda\|^2$ be a function on $\Lambda^2 n^* \otimes n$ and use Theorem 1.1 and Lemma 3.1.

It is possible that the above system of equations is not solvable, hence we assume that for some $k \in \mathbb{N}$ and every $\phi \in GL(k)$ there exists $\mu \in \mathfrak{N}_k$ such that $\phi.\mu$ isn't a Ricci soliton. \square

Notation 3.3. *Suppose that $DGL(k) := \{\phi \in GL(k); \phi \text{ is diagonal}\}$, For any $\mu \in \mathfrak{N}_k \subseteq \Lambda^2 n^* \otimes n$, set*

$$\mu(X_i, X_j) = \sum_v c_{ijv} X_v, \quad \phi = \text{diag}(\phi_1, \phi_2, \dots, \phi_k),$$

then

$$(\phi.\mu)(X_i, X_j) = \sum_{x_{ijv} \neq 0} x_{ijv} X_v; \quad x_{ijv} = \frac{\phi_v c_{ijv}}{\phi_i \phi_j}$$

and for any i, j and v such that $\langle \mu(X_i, X_j), X_v \rangle \neq 0$, the diagonal elements of $R_{\phi, \mu}$ are equal to

$$\begin{aligned}(R_{\phi, \mu})_{ii} &= -\frac{1}{2}(-x_{ijv}^2 + \sum_{\substack{rst \\ x_{rst} \neq x_{ijv}}} \delta^{rst, i} x_{rst}^2), \\ (R_{\phi, \mu})_{jj} &= -\frac{1}{2}(-x_{ijv}^2 + \sum_{\substack{rst \\ x_{rst} \neq x_{ijv}}} \delta^{rst, j} x_{rst}^2), \\ (R_{\phi, \mu})_{vv} &= \frac{1}{2}(x_{ijv}^2 + \sum_{\substack{rst \\ x_{rst} \neq x_{ijv}}} \delta^{rst, v} x_{rst}^2),\end{aligned}$$

where $\delta^{rst, i}$, $\delta^{rst, j}$ and $\delta^{rst, v}$ are equal to 0, 1 or -1.

Lemma 3.4. [6]. Let $(n, \mu, \langle \cdot, \cdot \rangle)$ be a Lie algebra and $P_\mu = \text{Sym}(n) \cap \text{Der}(n)$, then $R_\mu \perp P_\mu$ with inner product $\text{tr}(AB)$ on $\text{Sym}(n) \times \text{Sym}(n)$.

Theorem 3.5. (Lauret theory) For every $\phi \in GL(k)$ and $\mu \in \mathfrak{N}_k$ if $\phi \cdot \mu \in S_r$ and R_μ is diagonal, then for any i, j and v such that $\langle \mu(X_i, X_j), X_v \rangle \neq 0$, the following statements are equivalent:

- (i) $R_{\psi \cdot \mu} = c_{\psi \cdot \mu} I + D_{\psi \cdot \mu}$; $D_{\psi \cdot \mu} \in \text{Der}(\psi \cdot \mu)$
- (ii) $c_{\psi \cdot \mu}$ and a_{ijv}^2 's are solutions of the system

$$\begin{cases} \sum_{i, j, v} x_{ijv}^2 = r^2 \\ \frac{\partial F_k(\phi \cdot \mu)}{\partial u_{ijv}} \Big|_{x_{ijv} := a_{ijv}} = -c_{\psi \cdot \mu} \end{cases}$$

where $(\phi \cdot \mu)(X_i, X_j) = \sum_{x_{ijv} \neq 0} x_{ijv} X_v$, $u_{ijv} = x_{ijv}^2$, $\psi \cdot \mu = \phi \cdot \mu|_{x_{ijv} := a_{ijv}}$.

Proof. By Lemma 3.2 and the chain rule, it is easy to see that $\psi \cdot \mu$ is a Ricci soliton if and only if variation t and a_{ijv}^2 's are solutions of the system

$$\begin{cases} \sum_{i, j, v} x_{ijv}^2 = r^2 \\ \frac{\partial F_k(\phi \cdot \mu)}{\partial u_{ijv}} \Big|_{x_{ijv} := a_{ijv}} = t \end{cases}$$

Now we shall obtain the Lagrangian coefficient. By Lemma 3.3 it is easy to see that $\text{tr} R_\mu^2 = c_\mu \text{tr} R_\mu$. Also $\text{tr} R_\mu = -\frac{1}{2} \|\mu\|^2$. Thus $F_k = -c_\mu r^2$. Consequently $\frac{\partial F_k(\phi \cdot \mu)}{\partial u_{ijv}} \Big|_{x_{ijv} := a_{ijv}} = -c_\mu$. \square

Finally, we exhibit a rank-two Einstein solvmanifold of dimension 8 and a rank-three Einstein solvmanifold of dimension 15.

Example 3.6. Let $\mu = \mu_1 \oplus \mu_2$, where $\mu_1(X_1, X_2) = X_5$, $\mu_2(X_3, X_4) = X_6$ and $\phi = \text{diag}(\phi_1, \phi_2, \dots, \phi_6)$, then

$$\phi \cdot \mu_1(X_1, X_2) = \frac{\phi_5}{\phi_1 \phi_2} X_5 \quad \text{and} \quad \phi \cdot \mu_2(X_3, X_4) = \frac{\phi_6}{\phi_3 \phi_4} X_6.$$

Let $x := \frac{\phi_5}{\phi_1\phi_2}$, $y := \frac{\phi_6}{\phi_3\phi_4}$. Using Theorem 3.1, it is easy to see $x^2 = y^2 = 1$. Therefore $D_{\mu_1} = \text{diag}(1, 1, 0, 0, 2, 0)$ and $D_{\mu_2} = \text{diag}(0, 0, 1, 1, 0, 2)$. If $(n_1, \mu_1, \langle \cdot, \cdot \rangle_1)$ and $(n_2, \mu_2, \langle \cdot, \cdot \rangle_2)$ are nilradical Lie algebras, define the Lie algebra s with simply connected Lie group S using the following direct sum

$$s = \mathbb{R}Z_1 \oplus n_1 \oplus \mathbb{R}Z_2 \oplus n_2$$

endowed with the inner product $\langle \cdot, \cdot \rangle$ defined by

$$\langle Z_i, Z_j \rangle = 4\delta_{ij}, \langle Z_i, n \rangle = 0, \langle \cdot, \cdot \rangle|_{n_i \times n_i} = \langle \cdot, \cdot \rangle_i; \quad 1 \leq i, j \leq 2,$$

where the Lie bracket $[\cdot, \cdot]$ on s is defined by

$$\begin{aligned} [Z_1, X_1] &= -[X_1, Z_1] = X_1, & [Z_1, X_2] &= -[X_2, Z_1] = X_2, \\ [Z_1, X_5] &= -[X_5, Z_1] = 2X_5, & [Z_2, X_3] &= -[X_3, Z_2] = X_3, \\ [Z_2, X_4] &= -[X_4, Z_2] = X_4, & [Z_2, X_6] &= -[X_6, Z_2] = 2X_6, \\ [Z_i, Z_j] &= 0, & [\cdot, \cdot]_s|_{n_i \times n_i} &= \mu_i; \quad 1 \leq i, j \leq 2 \end{aligned}$$

and it is equal to zero otherwise. It is easy to see that $R_{[\cdot, \cdot]_s} = -\frac{3}{2}I_{8 \times 8}$ which implies that S_μ is Einstein of rank 2.

Example 3.7. Let $\mu = \mu_1 \oplus \mu_2 \oplus \mu_3$, where μ_1 and μ_2 are Lie algebras given in Example 1 and $\mu_3 \in \mathfrak{N}_6$ is given by

$$\mu_3(X_7, X_i) = X_{i+1}; \quad 8 \leq i \leq 11.$$

Every $\phi \cdot \mu_3 \in DGL(6) \cdot \mu_3 \cap S_{\sqrt{\frac{30}{13}}}$ is equal to

$$\phi \cdot \mu_3(X_7, X_i) = a_{7,i,i+1} X_{i+1}; \quad i = 8, 9, 10, 11.$$

By Lauret theory it is easy to see that a critical point of F_6 restricted to the leaf $\sum_{7 < i < 12} a_{7,i,i+1}^2 = \frac{20}{13}$ is equal to $\mu_3\{a_{7,i,i+1}\}$ where

$$a_{7,8,9}^2 = \frac{12}{13}, \quad a_{7,9,10}^2 = \frac{9}{13}, \quad a_{7,10,11}^2 = \frac{3}{13}, \quad a_{7,11,12}^2 = \frac{9}{13}, \quad c_{\psi \cdot \mu_3} = -\frac{3}{2}$$

and

$$D_{\mu_3(a_{7,i,i+1})} = \text{diag}(0, 0, 0, 0, 0, 0, \frac{15}{26}, \frac{42}{26}, \frac{36}{26}, \frac{51}{26}, \frac{48}{26}, \frac{51}{26}).$$

Let $(n_3, \mu_3(a_{7,i,i+1}), \langle \cdot, \cdot \rangle_3)$ be a Ricci soliton. Define Lie algebra s with simply connected Lie group S using the following direct sum

$$s = \mathbb{R}Z_1 \oplus n_1 \oplus \mathbb{R}Z_2 \oplus n_2 \oplus \mathbb{R}Z_3 \oplus n_3$$

endowed with the inner product $\langle \cdot, \cdot \rangle$ which is defined by

$$\langle Z_1, Z_1 \rangle = 4, \langle Z_2, Z_2 \rangle = 4, \langle Z_3, Z_3 \rangle = \frac{243}{13}, \langle \cdot, \cdot \rangle|_{n_i \times n_i} = \langle \cdot, \cdot \rangle_i$$

and it is equal to zero otherwise. Lie bracket $[\cdot, \cdot]$ on S defined by

$$\begin{aligned} [Z_1, X_1] &= -[X_1, Z_1] = X_1, & [Z_1, X_2] &= -[X_2, Z_1] = X_2, \\ [Z_1, X_5] &= -[X_5, Z_1] = 2X_5, & [Z_2, X_3] &= -[X_3, Z_2] = X_3, \\ [Z_2, X_4] &= -[X_4, Z_2] = X_4, & [Z_2, X_6] &= -[X_6, Z_2] = 2X_6, \end{aligned}$$

$$\begin{aligned}
[Z_3, X_7] &= -[X_7, Z_3] = \frac{15}{26}X_7, & [Z_3, X_8] &= -[X_8, Z_3] = \frac{42}{26}X_8, \\
[Z_3, X_9] &= -[X_9, Z_3] = \frac{36}{26}X_9, & [Z_3, X_{10}] &= -[X_{10}, Z_3] = \frac{51}{26}X_{10}, \\
[Z_3, X_{11}] &= -[X_{11}, Z_3] = \frac{48}{26}X_{11}, & [Z_3, X_{12}] &= -[X_{12}, Z_3] = \frac{51}{26}X_{12}, \\
[\cdot, \cdot]_{n_i \times n_i} &= \mu_{n_i}; \quad i = 1, 2
\end{aligned}$$

and otherwise is equal to zero. It is easy to check that $R_{[\cdot, \cdot]} = -\frac{3}{2}I_{15 \times 15}$ which implies that S_μ is Einstein of rank 3.

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REFERENCES

1. A. Besse, *Einstein manifolds*, *Ergeb. Math.* 10, 1987.
2. J. Heber, *Non-compact homogeneous Einstein spaces*, *Invent. Math.*, **133**, (1998), 279-352.
3. S. Helgason, *Differential Geometry Lie Groups and Symmetric Spaces*, Academic Press, New York, 1978.
4. S. Hong and M. M. Tripathi, Ricci curvature of submanifolds of a Sasakian space form, *Iranian Journal of Mathematical Sciences and Informatics*, **1**, (2006), 31-51.
5. J. Lauret, *Finding Einstein solvmanifolds by a variational method*, *Math. Z.*, **241**, (2003), 83-99.
6. J. Lauret, *Ricci soliton homogeneous nilmanifolds*, *Math. Ann.*, **319**, (2001), 715-733.
7. A. M. Tripathi, N. Mathur and S. Srivastava, A study of Nilpotent groups through right transversals, *Iranian Journal of Mathematical Sciences and Informatics*, **4**, (2009), 49-54.
8. J. Lauret, *Standard Einstein solvmanifolds as critical points*, *Quart. J. Math.*, **52**, (2001), 463-470.
9. C. Will, *Rank-one Einstein solvmanifolds of dimension 7*, *Differential Geom. Appl.*, **19**, (2003), 307-318.