

A QUASILINEAR PARABOLIC EQUATION WITH INHOMOGENEOUS DENSITY AND ABSORPTION

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ABSTRACT. We deal with the initial-boundary value problem for a quasilinear degenerate parabolic equation with inhomogeneous density and absorption, which appears in a number of applications to describe the evolution of diffusion processes, in particular non-Newtonian flow in a porous medium. We discuss the extinction of solution and the finite speed of propagation of perturbations.

1. Introduction

We consider the following equation,

$$(1.1) \quad \rho(x) \frac{\partial u}{\partial t} = \operatorname{div}(|\nabla u|^{p-2} \nabla u) - u^q, \quad (x, t) \in Q,$$

with the initial-boundary conditions,

$$(1.2) \quad u(x, 0) = u_0(x),$$

$$(1.3) \quad u(x, t) = 0, \quad x \in \partial\Omega,$$

where, $Q = \Omega \times (0, \infty)$, $\Omega \subset R^n$ is a bounded smooth domain, $p > 1$, $q \geq p - 1$, $u_0(x) \in C(\bar{\Omega}) \cap W_0^{1,p}(\Omega)$ is a nonzero nonnegative function and $\rho(x)$ denotes the density. We prefer to consider a typical case of

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$\rho(x)$, that is, $\rho(x) = (1 + |x|)^{-l}, l > 0$ [11]. The equation (1.1) is a prototype of a certain class of degenerate equations and appears to be relevant to the theory of non-Newtonian fluids [1]. For the case $\rho(x) = 1$, there have been many results about the existence, uniqueness and the regularity of the solutions. We refer the readers to the bibliography given in [4, 9, 13] and the references therein. Eidus [5], and Eidus and Kamin [6] considered the following problem:

$$\begin{aligned}\rho(x)u_t &= \Delta G(u), (x, t) \in Q_T = \Omega \times (0, T), \\ u(x, 0) &= u_0(x), \\ u|_{\partial\Omega} &= 0.\end{aligned}$$

They proved that the problem had a unique nonnegative solution, satisfying the condition,

$$\lim_{R \rightarrow \infty} R^{1-n} \int_{S(R)} \int_0^T G(u(x, t)) dt dx = 0,$$

where, $S(R) = \{x : x \in \Omega, |x| = R\}$.

Recently, Tedeev [10] considered the equation

$$\rho(x) \frac{\partial u}{\partial t} = \operatorname{div}(u^{m-1} |Du|^{\lambda-1} Du),$$

where, $\lambda > 0, m + \lambda - 2 > 0$ and with $\rho(x)$ being a positive continuous function. They examined under which conditions on $\rho(x)$, the corresponding nonnegative solutions of the Cauchy problems possessed the finite speed of propagations or the interface blow-up phenomena.

The equation (1.1) is degenerate, if $p > 2$, or singular, if $1 < p < 2$. Therefore, problem (1.1)-(1.3) does not admit classical solutions, in general. So, we study weak solutions in the sense of the following definition.

Definition 1.1. A nonnegative function u is said to be a weak solution of the problem (1.1)-(1.3), if u satisfies following conditions:

$$(1.4) \quad \begin{aligned}u &\in C(\bar{\Omega} \times (0, +\infty)) \cap L^\infty(0, \infty; W_0^{1,p}(\Omega)), \\ \int_0^T \int_\Omega \left(\rho(x) u \frac{\partial \varphi}{\partial t} - |\nabla u|^{p-2} \nabla u \nabla \varphi - u^q \varphi \right) dx dt &= 0,\end{aligned}$$

for all $\varphi \in C_0^\infty(Q_T)$, and $T \in (0, +\infty)$.

The existence proof for problem (1.1)-(1.3) is similar to the one for the case $\rho(x) = 1$. Here, our interest is to investigate the extinction

of solutions and the finite speed of propagation of perturbations. As is well known, one important property of solutions of the porous medium equation is the finite speed of propagation of perturbations. So, from the point of view of physical background, it seems to be natural to investigate this property for the equation (1.1). On the other hand, the mathematical description of this property is that if $\text{supp } u_0$ is bounded, then for any $t > 0$, $\text{supp } u(\cdot, t)$ is also bounded. So, from a mathematical viewpoint, this problem seems to be quite interesting. The monotonicity of support of weak solutions for the p -Laplacian equation was obtained by Yuan[12]. To prove the extinction of solution, here we use some ideas in [12]. We first construct a supersolution, and then use the comparison principle. Our method is different from the one given in [5] for the proof of the finite speed of propagation. We adopt the Bernis energy approach (see [2], [3]) and the main technical tools are weighted Nirenberg's inequality and Hardy's inequality.

2. Comparison principle

Here, we prove some lemmas.

Lemma 2.1. For $\varphi \in L^\infty(t_1, t_2; W_0^{1,p}(\Omega))$ with $\varphi_t \in L^2((t_1, t_2) \times \Omega)$, the weak solutions u of the problem (1.1)-(1.3) on Q_T satisfies

$$\begin{aligned} & \int_{\Omega} \rho(x)u(x, t_1)\varphi(x, t_1)dx + \int_{t_1}^{t_2} \int_{\Omega} \left(\rho(x)u \frac{\partial \varphi}{\partial t} - |\nabla u|^{p-2} \nabla u \nabla \varphi \right) dxdt \\ &= \int_{t_1}^{t_2} \int_{\Omega} u^q \varphi dxdt + \int_{\Omega} \rho(x)u(x, t_2)\varphi(x, t_2)dx. \end{aligned}$$

In particular, for $\varphi \in W_0^{1,p}(\Omega)$, we have

$$(2.1) \quad \begin{aligned} & \int_{\Omega} \rho(x)(u(x, t_1) - u(x, t_2))\varphi dx \\ & - \int_{t_1}^{t_2} \int_{\Omega} |\nabla u|^{p-2} \nabla u \nabla \varphi dxdt - \int_{t_1}^{t_2} \int_{\Omega} u^q \varphi dxdt = 0. \end{aligned}$$

Proof. From $\varphi \in L^\infty(t_1, t_2; W_0^{1,p}(\Omega))$ and $\varphi_t \in L^2((t_1, t_2) \times \Omega)$, it follows that there exists a sequence of functions $\{\varphi_k\}$, for fixed $t \in$

$(t_1, t_2), \varphi_k(\cdot, t) \in C_0^\infty(\Omega)$, and as $k \rightarrow \infty$,

$$\|\varphi_{kt} - \varphi_t\|_{L^2((t_1, t_2) \times \Omega)} \rightarrow 0, \quad \|\varphi_k - \varphi\|_{L^\infty(t_1, t_2; W_0^{1,p}(\Omega))} \rightarrow 0.$$

Choose a function $j(s) \in C_0^\infty(\mathbb{R})$ such that $j(s) \geq 0$, for $s \in \mathbb{R}$; $j(s) = 0$, $\forall |s| > 1$, and $\int_{\mathbb{R}} j(s) ds = 1$. For $h > 0$, define $j_h(s) = \frac{1}{h} j(\frac{s}{h})$ and

$$\eta_h(t) = \int_{t-t_2+2h}^{t-t_1-2h} j_h(s) ds.$$

Clearly, $\eta_h(t) \in C_0^\infty(t_1, t_2)$, and $\lim_{h \rightarrow 0^+} \eta_h(t) = 1$, for all $t \in (t_1, t_2)$. In the definition of weak solutions, choose $\varphi = \varphi_k(x, t)\eta_h(t)$. We have

$$\begin{aligned} & \int_{t_1}^{t_2} \int_{\Omega} \rho(x) u \varphi_k j_h(t - t_1 - 2h) dx dt - \int_{t_1}^{t_2} \int_{\Omega} |\nabla u|^{p-2} \nabla u \nabla \varphi_k \eta_h dx dt \\ & + \int_{t_1}^{t_2} \int_{\Omega} \rho(x) u \varphi_{kt} \eta_h dx dt - \int_{t_1}^{t_2} \int_{\Omega} \rho(x) u \varphi_k j_h(t - t_2 + 2h) dx dt \\ & = \int_{t_1}^{t_2} \int_{\Omega} u^q \varphi_k \eta_h dx dt. \end{aligned}$$

Observe that

$$\begin{aligned} & \left| \int_{t_1}^{t_2} \int_{\Omega} \rho(x) u \varphi_k j_h(t - t_1 - 2h) dx dt - \int_{\Omega} (\rho(x) u \varphi_k)|_{t=t_1} dx \right| \\ & = \left| \int_{t_1+h}^{t_1+3h} \int_{\Omega} \rho(x) u \varphi_k j_h(t - t_1 - 2h) dx dt \right. \\ & \quad \left. - \int_{t_1+h}^{t_1+3h} \int_{\Omega} (\rho(x) u \varphi_k)|_{t=t_1} j_h(t - t_1 - 2h) dx dt \right| \\ & \leq \sup_{t_1+h < t < t_1+3h} \int_{\Omega} |(\rho(x) u \varphi_k)|_t - (\rho(x) u \varphi_k)|_{t_1}| dx, \end{aligned}$$

and $u \in C(Q)$. We see that the right hand side tends to zero, as $h \rightarrow 0$. Similarly,

$$\left| \int_{t_1}^{t_2} \int_{\Omega} \rho(x) u \varphi_k j_h(t - t_2 + 2h) dx dt - \int_{\Omega} (\rho(x) u \varphi_k)|_{t=t_2} dx \right| \rightarrow 0.$$

Letting $h \rightarrow 0$ and $k \rightarrow \infty$, we obtain:

$$\begin{aligned} & \int_{\Omega} \rho(x) u(x, t_1) \varphi(x, t_1) dx + \int_{t_1}^{t_2} \int_{\Omega} \left(\rho(x) u \frac{\partial \varphi}{\partial t} - |\nabla u|^{p-2} \nabla u \nabla \varphi \right) dx dt \\ & = \int_{t_1}^{t_2} \int_{\Omega} u^q \varphi dx dt + \int_{\Omega} \rho(x) u(x, t_2) \varphi(x, t_2) dx. \end{aligned}$$

In particular, for $\varphi \in W_0^{1,p}(\Omega)$, we have

$$\int_{\Omega} \rho(x)(u(x, t_1) - u(x, t_2))\varphi dx - \int_{t_1}^{t_2} \int_{\Omega} |\nabla u|^{p-2} \nabla u \nabla \varphi dx dt - \int_{t_1}^{t_2} \int_{\Omega} u^q \varphi dx dt = 0,$$

which completes the proof. \square

For a fixed $\tau \in (0, T)$, set h satisfying $0 < \tau < \tau + h < T$. Let $t_1 = \tau$, $t_2 = \tau + h$, and then multiply (2.1) by $\frac{1}{h}$, for $\varphi \in W_0^{1,p}(\Omega)$, to obtain:

$$(2.2) \quad \int_{\Omega} \rho(x)(u_h(x, \tau))_{\tau} \varphi dx + \int_{\Omega} (|\nabla u|^{p-2} \nabla u)_h \nabla \varphi dx + \int_{\Omega} u_h^q \varphi dx = 0,$$

where,

$$u_h(x, t) = \begin{cases} \frac{1}{h} \int_t^{t+h} u(\cdot, \tau) d\tau, & t \in (0, T-h), \\ 0, & t > T-h. \end{cases}$$

Lemma 2.2. (Comparison principle) *Let u be a weak solution of (1.1)-(1.3). If v satisfies*

$$\rho(x) \frac{\partial v}{\partial t} \geq \operatorname{div} (|\nabla v|^{p-2} \nabla v) - v^q,$$

in the sense of distributions, and

$$v(x, 0) \geq u(x, 0),$$

$$v(x, t) \geq u(x, t), \quad x \in \partial\Omega,$$

then we have

$$v(x, t) \geq u(x, t), \quad \text{for all } (x, t) \in Q.$$

Proof. By (2.2), we have for $\varphi \in W_0^{1,p}(\Omega)$,

$$\int_{\Omega} \rho(x)(u(x, \tau) - v(x, \tau))_{h\tau} \varphi(x) dx + \int_{\Omega} (u^q - v^q)_h(x, \tau) \varphi dx + \int_{\Omega} (|\nabla u|^{p-2} \nabla u - |\nabla v|^{p-2} \nabla v)_h(x, \tau) \nabla \varphi dx \leq 0.$$

For a fixed τ , we take $\varphi(x) = [(u - v)_h]_+$. By the property of the Steklov mean value and noting that $v(x, t) \geq u(x, t), x \in \partial\Omega$, we see

that $\varphi(x) = [(u - v)_h]_+ \in W_0^{1,p}(\Omega)$. Substituting this function into the above integral equality, we obtain:

$$\begin{aligned} & \int_{\Omega} \rho(x)(u(x, \tau) - v(x, \tau))_{h\tau} [(u - v)_h]_+ dx \\ & \leq - \int_{\Omega} [(|\nabla u|^{p-2} \nabla u - |\nabla v|^{p-2} \nabla v)_h](x, \tau) \nabla [(u - v)_h]_+ dx \\ & \quad - \int_{\Omega} [(u^q - v^q)_h](x, \tau) [(u - v)_h]_+ dx. \end{aligned}$$

Integrating the above equality with respect to τ over $(0, t)$, we have

$$\begin{aligned} & \int_{\Omega} \rho(x) [(u - v)_h]_+^2(x, t) dx - \int_{\Omega} \rho(x) [(u - v)_h]_+^2(x, 0) dx \\ & \leq - \int_{\Omega} [(|\nabla u|^{p-2} \nabla u - |\nabla v|^{p-2} \nabla v)_h](x, \tau) \nabla [(u - v)_h]_+ dx \\ & \quad - \int_{\Omega} [(u^q - v^q)_h](x, \tau) [(u - v)_h]_+ dx, \end{aligned}$$

It is easily seen that

$$\lim_{h \rightarrow 0} \int_{\Omega} \rho(x) [(u - v)_h]_+(x, 0) dx = 0.$$

Letting $h \rightarrow 0$, we have

$$\int_{\Omega} \rho(x) |(u - v)_+|^2(x, t) dx \leq 0,$$

that is, $\int_{\Omega} |(u - v)_+|^2 dx = 0$. Therefore, $v \geq u$, and the proof is complete. \square

3. Extinction and monotonicity of support

We now to prove the following theorem.

Theorem 3.1. *Let u be a nonnegative weak solution of the problem (1.1)-(1.3), and $p > 2$. Then,*

$$\text{supp}(\cdot, s) \subset \text{supp}(\cdot, t),$$

for all s, t with $0 < s < t$.

Lemma 3.2. *Let u be a nonnegative weak solution of the problem (1.1)-(1.3) If $p > 2$, then*

$$\frac{\partial u}{\partial t} \geq -\frac{u}{(p-2)t},$$

in the sense of distributions.

Proof. Denote:

$$u_r(x, t) = ru(x, r^{p-2}t), \text{ for all } (x, t) \in Q, \quad r > 1.$$

By $p - 1 \leq q$ and $r > 1$, we have

$$\rho(x) \frac{\partial u_r}{\partial t} \geq \operatorname{div}(|\nabla u_r|^{p-2} \nabla u_r) - u_r^q,$$

$$(3.1) \quad u_r(x, 0) = ru_0(x),$$

$$(3.2) \quad u_r(x, t) = 0, \quad x \in \partial\Omega.$$

Noting that $r > 1$, and using (1.2), (3.1), and (3.2), we get

$$(3.3) \quad u_r(x, 0) \geq u_0(x),$$

$$(3.4) \quad u_r(x, t) = u(x, t), \quad x \in \partial\Omega.$$

Applying the comparison principle, we have

$$(3.5) \quad u_r(x, t) \geq u(x, t).$$

For $p > 2$, by (3.5), we obtain:

$$\frac{[u(x, \lambda t)]^{p-2} - [u(x, t)]^{p-2}}{\lambda t - t} \geq \frac{(1/\lambda - 1)[u(x, t)]^{p-2}}{\lambda t - t},$$

where, $\lambda = r^{p-2}$. Letting $\lambda \rightarrow 1^+$, we get

$$\frac{\partial}{\partial t} [u(x, t)]^{p-2} \geq -\frac{1}{t} [u(x, t)]^{p-2},$$

in the distribution, which implies that lemma holds. Thus the proof is now complete. \square

Proof of Theorem 3.1. For $p > 2$, from Lemma 3.2 we obtain:

$$\frac{\partial (t^{1/(p-2)}u)}{\partial t} \geq 0.$$

\square

Theorem 3.3. *Let u be a weak solution of (1.1)-(1.3). If $1 < p < 2$, then there exists a time T such that*

$$u(x, t) = 0,$$

for all $(x, t) \in \Omega \times (T, \infty)$.

Proof. Denote $s_+ = \max\{s, 0\}$, for all $s \in (-\infty, +\infty)$.

Define an auxiliary function,

$$(3.6) \quad v(x, t) = k(T-t)_+^{1/(2-p)} \ln(m + x_1 + \cdots + x_n),$$

where,

$$(3.7) \quad k = \left[\frac{(p-1)(2-p)n^{p/2}}{(2m)^p \ln(2m)} \right]^{1/(2-p)}, \quad T = \left(\frac{\max \|u_0\|_\infty}{k \ln 2} \right)^{2-p},$$

and

$$(3.8) \quad m = \sup_{x \in \Omega} \{|x_1| + \cdots + |x_n|\} + 2.$$

Compute

$$(3.9) \quad \frac{\partial v}{\partial t} = -\frac{k}{2-p} (T-t)_+^{(p-1)/(2-p)} \ln(m + x_1 + \cdots + x_n),$$

$$(3.10) \quad \begin{aligned} & \operatorname{div} \left(|\nabla v|^{p-2} \nabla v \right) \\ &= \operatorname{div} \left\{ k^{p-1} (T-t)_+^{(p-1)/(2-p)} n^{(p-2)/2} \right. \\ & \quad \left. \left((m + x_1 + \cdots + x_n)^{1-p}, \dots, (m + x_1 + \cdots + x_n)^{1-p} \right) \right\} \\ &= - (p-1) k^{p-1} (T-t)_+^{(p-1)/(2-p)} \\ & \quad \cdot n^{(p-2)/2} (m + x_1 + \cdots + x_n)^{-p}, \end{aligned}$$

and

$$v^q = k^q (T-t)_+^{q/(2-p)} \ln^q(m + x_1 + \cdots + x_n).$$

By $\rho(x) = (1 + |x|)^{-l}$, $l > 0$, and using (3.6)-(3.10), we get

$$\rho(x) \frac{\partial v}{\partial t} \geq \frac{\partial}{\partial x} \left(\left| \frac{\partial v}{\partial x} \right|^{p-2} \frac{\partial v}{\partial x} \right) - v^q,$$

and

$$\begin{aligned} v(x, 0) &\geq u(x, 0), \\ v(x, t) &\geq u(x, t) = 0, \quad x \in \partial\Omega. \end{aligned}$$

Applying Lemma 2.2, we obtain:

$$u(x, t) \leq v(x, t),$$

for all $(x, t) \in Q$. By the definition of $v(x, t)$, we have

$$u(x, t) \leq v(x, t) = 0,$$

for all $(x, t) \in \Omega \times (T, +\infty)$. Thus, the proof is complete. \square

4. Finite speed of propagation of perturbations

Here, we study the property of finite speed of perturbations.

Theorem 4.1. *Assume $p > 2$, $|\xi_n(0)| \leq a$, $|\sigma_n(0)| \leq b$, and u is the weak solution of the problem (1.1)-(1.3). Then, for any fixed $t > 0$, we have*

$$\begin{aligned} \sigma_n(t) - \sigma_n(0) &\leq C_1 t^\mu \left(\int_0^T \int_\Omega |\nabla u|^p dx dt \right)^\beta, \\ \xi_n(t) - \xi_n(0) &\geq -C_2 t^\mu \left(\int_0^T \int_\Omega |\nabla u|^p dx dt \right)^\beta, \end{aligned}$$

where, C_1, C_2 are constants depending on $n, p, a, b, u_0(x)$, $\sigma_n(t) = \sup\{z; x \in \text{suppu}(\cdot, t)\}$, $\xi_n(t) = \inf\{z; x \in \text{suppu}(\cdot, t)\}$, $z = x_n$, $\beta > 0, \mu > 0$, and $a, b > 0$ are constants independent of t .

Proof. First, we discuss the following Dirichlet problem,

$$(4.1) \quad \rho(x) \frac{\partial u}{\partial t} = \text{div}((|\nabla u|^2 + \varepsilon)^{\frac{p-2}{2}} \nabla u) - u^q,$$

$$(4.2) \quad u(x, t) = 0, \quad x \in \partial\Omega,$$

$$(4.3) \quad u(x, 0) = u_{0\varepsilon}(x).$$

It is well known that (4.1)-(4.3) has a nonnegative classical solution u_ε [8].

Multiplying (4.1) by $(z - y)_+^s u_\varepsilon(x)$, $s \geq 2p$, $y \geq \sigma_n(0)$, and letting $\varepsilon \rightarrow 0$, we have,

$$\begin{aligned} I &= \frac{1}{2} \int_\Omega \rho(x) (z - y)_+^s |u(x, t)|^2 dx + \int_0^t \int_\Omega (z - y)_+^s u^{q+1} dx d\tau \\ &= - \int_0^t \int_\Omega |\nabla u|^{p-2} \nabla u \nabla((z - y)_+^s u) dx d\tau. \end{aligned}$$

Then, we have,

$$\begin{aligned} I &= - \int_0^t \int_{\Omega} |\nabla u|^{p-2} \nabla u \nabla((z-y)_+^s u) dx d\tau \\ &= - \int_0^t \int_{\Omega} (z-y)_+^s |\nabla u|^p dx d\tau \\ &\quad - \int_0^t \int_{\Omega} \nabla[(z-y)_+^s] u |\nabla u|^{p-2} \nabla u dx d\tau. \end{aligned}$$

Using Hölder inequality,

$$\begin{aligned} I &\leq - \int_0^t \int_{\Omega} (z-y)_+^s |\nabla u|^p dx d\tau + \frac{1}{2} \int_0^t \int_{\Omega} (x-y)_+^s |\nabla u|^p dx d\tau \\ &\quad + C_1 \int_0^t \int_{\Omega} (z-y)_+^{s-p} |u|^p dx d\tau \\ &\leq - \frac{1}{2} \int_0^t \int_{\Omega} (x-y)_+^s |\nabla u|^p dx d\tau + C_1 \int_0^t \int_{\Omega} (z-y)_+^{s-p} |u|^p dx d\tau. \end{aligned}$$

Applying Hardy inequality [7], we obtain:

$$\int_{\Omega} (z-y)_+^{s-p} |u|^p dx \leq \left(\frac{p}{s-p+1} \right)^p \int_{\Omega} (z-y)_+^s |D_z u|^p dx.$$

Hence,

$$\begin{aligned} &\frac{1}{2} \int_{\Omega} \rho(x) (z-y)_+^s |u|^2 dx + \frac{1}{2} \int_0^t \int_{\Omega} (z-y)_+^s |\nabla u|^p dx d\tau \\ &\leq C \int_0^t \int_{\Omega} (z-y)_+^{s-p} |u|^p dx d\tau. \end{aligned}$$

Thus,

$$(4.4) \quad \sup_{0 < \tau \leq t} \int_{\Omega} \rho(x) (z-y)_+^s |u|^2 dx \leq C \iint_{Q_t} (z-y)_+^{s-p} |u|^p dx d\tau,$$

and

$$(4.5) \quad \iint_{Q_t} (z-y)_+^s |\nabla u|^p dx d\tau \leq C \iint_{Q_t} (z-y)_+^{s-p} |u|^p dx d\tau.$$

For (4.4), using Hardy inequality, again we have

$$(4.6) \quad \sup_{0 < \tau \leq t} \int_{\Omega} \rho(x) (z-y)_+^s |u|^2 dx \leq C \iint_{Q_t} (z-y)_+^s |\nabla u|^p dx d\tau.$$

Set

$$f_s(y) = \iint_{Q_t} (z-y)_+^s |\nabla u|^p dx d\tau, \quad f_0(y) = \int_0^t \int_{\{x \in \Omega, x_n > y\}} |\nabla u|^p dx d\tau.$$

From (4.5) and weighted Nirenberg inequality, we have

$$\begin{aligned} & f_{2p+1}(y) \\ & \leq C_1 \iint_{Q_t} (z-y)_+^{p+1} |u|^p dx d\tau \\ & \leq C \int_0^t \left(\int_{\Omega} (z-y)_+^{p+1} |\nabla u|^p dx \right)^a \left(\int_{\Omega} (z-y)_+^{p+1} |u|^2 dx \right)^{\frac{(1-a)p}{2}} d\tau, \end{aligned}$$

where, $\frac{1}{p} = a\left(\frac{1}{p} - \frac{1}{n+p+1}\right) + (1-a)\frac{1}{2}$. Therefore,

$$a = \frac{\frac{1}{p} - \frac{1}{2}}{\frac{1}{p} - \frac{1}{n+p+1} - \frac{1}{2}} < 1.$$

Using (4.6), we obtain:

$$\begin{aligned} & f_{2p+1}(y) \\ & \leq C \left(\iint_{Q_t} (z-y)_+^{p+1} |\nabla u|^p dx d\tau \right)^{\frac{(1-a)p}{2}} \int_0^t \left(\int_{\Omega} (z-y)_+^{p+1} |\nabla u|^p dx \right)^a d\tau \\ & \leq C [f_{p+1}(y)]^{(1-a)p/2} \left(\iint_{Q_t} (z-y)_+^{p+1} |\nabla u|^p dx d\tau \right)^a t^{1-a} \\ & \leq C f_{p+1}(y)^{(1-a)p/2+a} t^{1-a}. \end{aligned}$$

Denote $\lambda = 1 - a$ and $\mu = a + (1 - a)p/2$. Then, $\lambda > 0$ and $1 < \mu$. Applying Hölder's inequality, we have

$$\begin{aligned} & f_{2p+1}(y) \\ & \leq C t^\lambda \left[\iint_{Q_t} (z-y)_+^{p+1} |\nabla u|^p dx ds \right]^\mu \\ & \leq C t^\lambda \left[\iint_{Q_t} (z-y)_+^{2p+1} |\nabla u|^p dx ds \right]^{\frac{(p+1)\mu}{2p+1}} \left[\int_0^t \int_{\Omega} |\nabla u|^p dx ds \right]^{\frac{p\mu}{2p+1}} \\ & \leq C t^\lambda [f_{2p+1}(y)]^{\frac{(p+1)\mu}{2p+1}} [f_0(y)]^{\frac{p\mu}{2p+1}}. \end{aligned}$$

Therefore,

$$f_{2p+1}(y) \leq Ct^\lambda [f_0(y)]^{\frac{p\mu}{(2p+1)\sigma}}, \quad \sigma = 1 - \frac{p+1}{2p+1}\mu > 0.$$

Using Hölder's inequality again, we get

$$f_1(y) \leq [f_0(y)]^{\frac{2p}{2p+1}} [f_{2p+1}(y)]^{\frac{1}{2p+1}} \leq Ct^\gamma [f_0(y)]^{1+\theta},$$

where,

$$\gamma = \frac{\lambda}{(2p+1)\sigma}, \quad \theta = \frac{p\mu}{(2p+1)^2\sigma} - \frac{1}{2p+1} > 0.$$

Noticing that $f_1'(y) = -f_0(y)$, we obtain:

$$f_1'(y) \leq -Ct^{-\gamma/(\theta+1)} [f_1(y)]^{1/(\theta+1)}.$$

If $f_1(\sigma_n(0)) = 0$, then $\sigma_n(t) \leq b$. If $f_1(\sigma_n(0)) > 0$, then there exists a maximal interval $(\sigma_n(0), x^*)$, in which $f_1(y) > 0$, and

$$\left[f_1(y)^{\theta/(\theta+1)} \right]' = \frac{\theta}{\theta+1} \frac{f_1'(y)}{[f_1(y)]^{1/(\theta+1)}} \leq -Ct^{-\gamma/(\theta+1)}.$$

Integrating the above inequality over $(\sigma_n(0), x^*)$, we have

$$f_1(x^*)^{\theta/(\theta+1)} - f_1(\sigma_n(0))^{\theta/(\theta+1)} \leq -Ct^{-\gamma/(\theta+1)}(x^* - \sigma_n(0)),$$

which implies that

$$x^* \leq \sigma_n(0) + Ct^\gamma (f_0(y))^\theta,$$

noticing that $f_0(y)$ can be controlled by a constant C independent of y .

Similarly, we have

$$\xi_n(t) \geq \xi_n(0) - C_2 t^\mu \left(\int_0^T \int_\Omega |\nabla u|^p dx dt \right)^\beta.$$

The proof is now complete. \square

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