UNIQUENESS OF SOLUTION FOR A CLASS OF STEFAN PROBLEMS

K. Ivaz*

Faculty of Mathematics, Iran University of Science and Technology, Tehran 16844, Islamic Republic of Iran

Abstract

This paper deals with a theoretical mathematical analysis of one-dimensional solidification problem, in which kinetic undercooling is incorporated into the This temperature condition at the interface. A model problem with nonlinear kinetic law is considered. We prove a local result intimate for the uniqueness of solution of the corresponding free boundary problem.

Keywords: Free boundary problem; Nonlinear integral equation

Introduction

It is well known that in many industrial areas, the solidification process plays a significant role. Mathematical models of solidification including interface kinetics effects have been considered for quite some time (see [1], and references therein). This class of free boundary problems, which arises in a number of physical situations, is that of on equilibrium problems, in which the phase - change temperature is dependent on the velocity of the front at which the phase-change occurs (for more physical problems, see [3-7]). Here, we study a model problem with nonlinear kinetic law at the interface in the one-dimensional case. Specifically, let the curve with s(0)=b(0<b<1) be defined as the interface that separates the liquid and solid phases. With u denoting temperature (scaled so that is vanishes at equilibrium), we may write the system of equations as

$$u_t = K_l u_{xx}$$
 in $Q_1 = \{(x, t) | 0 < x < s(t), 0 < t \le T\},$ (1.1)

$$u_t = K_s u_{xx}$$
 in $Q_1 = \{(x,t) | 0 < x < s(t), 0 < t \le T\},$ (1.2)

and on the interface x = s(t) as

$$u_1 = u_2 = g_1(V(f)),$$
 (1.3)

$$Ku_{x}^{+} - Ku_{x}^{-} = g_{2}(V(t)),$$
 (1.4)

$$s(0) = b, \ 0 < b < 1,$$
 (1.5)

where K_l and K_s are thermal diffusivities of a liquid and a solid respectively, L > 0 is the latent heat and the superscripts + and - denote, respectively the right-hand and left-hand limits with respect to the special variable x. These equations are subject to the initial and boundary conditions

$$u(x,0) = \varphi_1(x) \quad 0 \le x \le b , \qquad (1.6)$$

$$u(x,0) = \varphi_2(x)$$
 $b \le x \le 1$, (1.7)

$$u(i-1,t) = f_i(t)$$
 $t \ge 0$, $(i=1,2)$ (1.8)

where

$$V(t) = \frac{ds(t)}{dt} \tag{1.9}$$

is the propagation velocity of the free boundary. The free boundary problem considered here was formulated in [1], where reduction the problem to an integral equation was given. In the context of solid fuel combustion, s(t) represents the boundary between the unburnt and burnt material, and u_1, u_2 , are the nondimensionalized temperature in the unburnt and burnt material respectively, (see [3-7] and references therein). The temperature at the free boundary controls its velocity $V(t) = g_1^{-1}(u_1(s(t),t))$. The heat exchange between the unburnt (x < s(t)) and burnt material is modeled by the boundary condition in (1.4) which, in principle, may be nonlinear.

Main Results

Theorem. Consider the problem (1.1)-(1.9). Suppose that the kinetic function and initial and boundary data satisfy the assumption (H_1) - (H_3) in [1]. Then the problem (1.1)-(1.9) has not more than one solution.

To prove uniqueness for $t < \sigma$ suppose that $u_0 = (u_{01}, u_{02}), s_0$ is another solution of (1.1)-(1.9) for $t < \sigma$ and $v_0(t) = (v_{01}(t), v_{02}(t))^T$ is another solution of integral equations (26) and (27) in [1]. It suffices to prove uniqueness, for any $\overline{\sigma} < \sigma$.

Let

$$\overline{M} = Max\{M, \underset{0 \leq t \leq \overline{\sigma}}{l.u.b} \big| v_0(t) \big| \}$$

where M introduced in section 4.2 in [1], and let be any positive number satisfying

$$\sqrt{\overline{\sigma}}$$
 <

$$\mathit{Min} \bigg\{ \Big(C_2 \Big\| \varphi_1^{\, \cdot} \Big\| + C_3 \Big\| f_1^{\, \cdot} \Big\| + C_4 + C_5 \overline{M} + C_6 + C_7 \Big)^{-1} \, \overline{\frac{M}{2}} \, ,$$

$$\left(D_{2}+D_{3}+D_{4}\overline{M}+D_{5}+D_{6}\left\|\varphi_{2}^{'}\right\|+D_{7}\left\|f_{2}^{'}\right\|\right)^{-1}\frac{\overline{M}}{2}\right\}$$

where the constants C_i and D_i , i=2,3,...,7 are simple combination of $\pi, b, \frac{1}{b}, M, M', M_2, M'_2, K$. Then by the same calculations in [1] which were used to prove that T maps $B_{M,\sigma}$ into itself (where T and $B_{M,\sigma}$ introduced in subsection 4.2 in [1]) and is a contraction one shows that T maps $B_{\overline{M}} = \overline{\sigma}$ into itself and is a contraction. Hence, there exists at most one fixed point of T in $B_{\overline{M},\overline{\sigma}}$. It follows that $v(t) = v_0(t)$ for $0 \le t \le \overline{\sigma}$, where v(t) is solution of integral equations (26) and (27) in [1]. Hence also $s(t) = s_0(t)$ $u(x,t) = u_0(x,t)$ if $0 \le t \le \overline{\sigma}$, $0 \le x \le s(t)$ and $s(t) \le x \le 1$. We next consider the system (1.1)-(1.9) for $t > \sigma$, i.e. (1.1)-(1.5), (1.8), (1.9) are considered for $t > \sigma$ (instead of $t \ge 0$) where as (1.6), (1.7) are replaced by $u_1(x, \overline{\sigma}) = u_1(x, \overline{\sigma})$ for $0 \le x \le s(\overline{\sigma})$, $u_2(x, \overline{\sigma}) = u_2(x, \overline{\sigma})$ for $s(\overline{\sigma}) < x < 1$.

This problem can again be transformed into integral equations (26), (27) in [1] extend to the present integral equation provided M is replaced by M_0 where

$$M_0 = \underline{l}.u.b$$
 $V(t)g_1(V(t))$

Similarly to section 4 in [1], we reduce the problem (1.1)-(1.9) for u_0 , s_0 in the interval $\sigma \le t < \sigma$ to an integral equation. Since $u_1(x,\sigma) = u_1(x,\sigma)$, $u_2(x,\sigma) = u_2(x,\sigma)$, the integral equation for v(t) and $v_0(t)$ coincide. Repeating now the same argument as before we conclude that for $v(t) = v_0(t)$ for any $\tilde{\sigma}$ satisfying

$$\begin{split} &\sqrt{\left(\widetilde{\sigma}-\overline{\sigma}\right)} < \\ &Min \bigg\{ \left(C_{2} \left\| \varphi_{1}^{\cdot} \right\| + C_{3} \left\| f_{1}^{\cdot} \right\| + C_{4} + C_{5} \overline{M_{0}} + C_{6} + C_{7} \right)^{-1} \frac{\overline{M_{0}}}{2}, \\ &\left(D_{2} + D_{3} + D_{4} \overline{M_{0}} + D_{5} + D_{6} \left\| \varphi_{2}^{\cdot} \right\| + D_{7} \left\| f_{2}^{\cdot} \right\| \right)^{-1} \frac{\overline{M_{0}}}{2} \bigg\} \end{split}$$

$$(\overline{M}_0 = Max\{M_0, \underset{\overline{\sigma} \le t \le \sigma}{l.u.b} |v_0(t)|\})$$

We can now proceed in the same manner as before in [2] step by step, nothing that in each step the time interval can be taken to be $\geq \varepsilon$ where satisfies

 $\sqrt{\varepsilon}$ <

$$Min \left\{ \left(C_{2} \left\| \varphi_{1}^{'} \right\| + C_{3} \left\| f_{1}^{'} \right\| + C_{4} + C_{5} \overline{M_{1}} + C_{6} + C_{7} \right)^{-1} \frac{\overline{M_{1}}}{2}, \right.$$

$$\left(D_{2}+D_{3}+D_{4}\overline{M_{1}}+D_{5}+D_{6}\left\| \varphi_{2}^{'}\right\| +D_{7}\left\| f_{2}^{'}\right\| \right)^{-1}\frac{\overline{M_{1}}}{2}$$

where

$$\overline{M}_1 = Max \{ \underset{\overline{\sigma} < t < \sigma}{l.u.b} \quad \left| V(t)g_1(V(t)) \right|, \underset{\overline{\sigma} < t < \sigma}{l.u.b} \quad \left| v_0(t) \right| \}$$

Having proved existence and uniqueness for all $t < \sigma$ where σ is any positive number satisfying (36) in [1]. Let us stress that the previous proof (see (38), (39) in [1]) shows also the following:

If instead of (1.1)-(1.9) for t>0 we consider (1.1)-(1.9) for $t>\lambda$, i.e., (1.1)-(1.5), (1.8), (1.9) hold for $t>\lambda$ and (1.6), (1.7) replaced by $u_1(x,\lambda)=u_1(x,\lambda)$ for $0< x< s(\lambda)$ and $u_2(x,\lambda)=u_2(x,\lambda)$ for $s(\lambda)< x\leq 1$ respectively, and if

$$|V(\lambda)g_1(V(\lambda))|, s(\lambda), \frac{1}{s(\lambda)}$$

are bounded independently of λ , then there exists a unique solution for the problem in an interval $\lambda \leq t \leq \lambda + \varepsilon$, where ε is some positive number independent of λ .

Since for any solution of (1.1)-(1.9) the function s(t) is monotone non-decreasing, $\frac{1}{s(\lambda)} \le \frac{1}{b}$. To complete

the proof of theorem it suffices to prove the following statement:

For every $t_0 > 0$ there exists an $\varepsilon > 0$ such that if the system (1.1)-(1.9) has a unique solution for all $t < t_0$, then it also has a unique solution for all $t < t_0 + \varepsilon$ in view of the previous remarks it suffices to show: If

u(x,t), s(t) is a solution of (1.1)-(1.9) for all $t < t_0$, then for all $\eta > 0$ sufficiently small, the functions

$$|l.u.b|V(t_0 - \eta)g_1(V(t_0 - \eta))|, s(t_0 - \eta)$$
 (2.1)

are bounded independently of η . If we prove that

$$|l.u.b|v(t)| < \infty,$$

then from (28) in [1] follows the boundedness of s(t) for $t < t_0$. Consequently, if we prove (2.2) then the proof of theorem is completed.

Proof of (2.2). We use for v(t) the integral equation which corresponds to the system (1.1)-(1.9) in the interval $t_0 - \mu < t < \mu$ (μ sufficiently small) in [1]. Since

$$u_1(0,t_0-\mu) = f_1(t_0-\mu), u_2(0,t_0-\mu) = f_2(t_0-\mu),$$

the equations are

$$v_1(t) = -g_1(V(t))V(t) +$$

$$2\int_{0}^{s(t_{0}-\mu)}u_{1\xi}(\xi,t_{0}-\mu)N(s(t),t;\xi,t_{0}-\mu)d\xi$$

$$+ \int_{t_0-\mu}^t v_1(\tau) G_x(s(t),t;s(\tau),\tau) d\tau$$

$$-2\!\int_{t_0}^t f_1'(\tau)N(s(t),t;0,\tau)d\tau +$$

$$2g_1(V(t_0-\mu))N(s(t),t;s(t_0-\mu),t_0-\mu)$$

$$+2\int_{t_0-\mu}^t g_1'(V(\tau))N(s(t),t;s(\tau),\tau)d\tau$$

$$-2\int_{t_{\tau}-t_{t}}^{t}g_{1}(V(\tau))V(\tau)G_{x}(s(t)-1,t;s(\tau),\tau)d\tau$$

$$\begin{split} v_2(t) &= -g_1(V(t))V(t) \\ &- 2g_1(V(t_0 - \mu))N(s(t) - 1, t; s(\tau) - 1, t_0 - \mu) \\ &+ \int_{t_0 - \mu}^t g_1'(V(\tau))N(s(t) - 1, t; s(\tau) - 1, \tau)d\tau \\ &- 2\int_{t_0 - \mu}^t v_2(t)G_x(s(t) - 1, t; s(\tau) - 1, \tau)d\tau \\ &+ 2\int_{s(t_0 - \mu)}^1 u_{2\xi}(\xi, t_0 - \mu)N(s(t) - 1, t; \xi, t_0 - \mu)d\xi \\ &+ \int_{t_0}^t f_2'(\tau)N(s(t) - 1, t; 0, \tau)d\tau \end{split}$$

In section 4 in [1] we proved $v_1(t)$ and $v_2(t)$ are bounded functions, we obtain that

$$\lim_{0 < t < t_0} |v(t)| < \infty$$

therefore we established theorem.

References

- Shidfar A. and Ivaz K. The Stefan problem with kinetic function at the free boundary. J. Sci. I. R. Iran, 10(2): 121-129 (1999).
- 2. Friedman A. *Partial Differential Equation of Parabolic Type*. Prentice-Hall (1964).
- 3. Frankel M. Free boundary problems and dynamical geometry associated with flames, In: Fife P., Linan A. and Williams F. (Eds.), *Dynamical Issues in Combustion Theory*. IMA Volumes in Mathematics and its Applications, **35**: 107-127 (1991).
- Shkadin K.G. and Mehrzhanov A.G. Propagation of a pulsating exothermic reaction front in the condensed phase. *Combust. Expl. Shock Waves*, 7: 15-22 (1971).
- Mehrzhanov A.G., Filonenko A.K. and Borovinskaya I.P. New phenomena in combustion of condensed systems. Dokl. Akad. Nauk USSR, 208: 892-894 (1973); Soviet Phys. Doki., 208: 122-125 (1973).
- Mehrzhanov A.G. SHS Prxesses: Combustion theory and practice. Arch. Combustiones, 1: 23-28 (1981).
- 7. Matkowsky B.J. and Sivashinsky G.I. Propagation of a pulsating reaction front in solid fuel combustion. *SIAM J. Appl. Math.* **35**: 230-255 (1978).